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## DEVICE, SYSTEM AND METHOD FOR MEASURING REICHENBACH CLOCK SYNCHRONIZATIONS

### 5 FIELD OF THE INVENTION

The present invention is directed to a device, system and method for measuring the Reichenbach clock synchronization coefficients and the resulting vector velocity of light.

### 10 BACKGROUND OF THE INVENTION

This invention relates in general to the determination of velocity as defined by the cosmic background Doppler shift by measuring the Reichenbach clock synchronization coefficients and the resulting vector velocity of light in the oscillations of photon tunneling times.

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Einstein first introduced the idea of an ultimate particle speed known as  $c$ , the speed of light, with the publication of his special theory of relativity in 1905. Since this publication, scientists have shown that  $c$  is not an upper-limit on a particle's speed, but a barrier to acceleration. These mathematical studies have shown that while it is not possible to accelerate an object to a velocity faster than light, it is possible for an object to have a velocity greater than  $c$ .

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In mathematical terms the one-way vacuum velocity of light from A to B is  $c(AB)$ , which is described by the equation:

$$c(AB) = c / 2\epsilon(AB) \quad (1)$$

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where  $\epsilon(AB)$  are the Reichenbach clock synchronization coefficients. The vector velocity  $c(AB)$  is only isotropic in a preferred reference frame and the round trip vacuum speed of light,  $c$ , is constant because:

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$$\epsilon(AB) = 1 - \epsilon(BA) \quad (2)$$

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In the Lorentz covariant theoretical work, without superluminal energy flow, any inertial frame can be chosen as the preferred frame. Only by using superluminal energy flow, then, can a single preferred reference frame be used to measure the vector vacuum velocity of light.

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The first evidence of energy moving at velocities greater than  $c$  was observed by radio engineers at the turn of the century. They learned that radio signals in the upper atmosphere traveled faster than light. The reason was that the radio waves were moving through ionized gas and not normal air. In effect, these radio waves' pulses have two different velocities, a group velocity, or the velocity of the pulse packet, and a phase velocity, the velocity of the individual waves within the group. In this example, the phase velocity of the radio waves, or the internal velocity of the individual waves within the radio wave pulse packet were moving faster than light. A more complete discussion of these early superluminal radio wave experiments can be found in the text *Faster Than Light*, by Nick Herbert, pg. 56-58, (1988).

Systems designed to transmit energy at superluminal velocities are also well-known in the art of quantum mechanics. One type of conventional superluminal energy transport method employs the phenomenon known as quantum barrier penetration, or tunneling. Under quantum theory, a quantum particle can be thought of as a wave packet, its width in space related to its velocity through the Heisenberg Uncertainty Relation. A common interpretation of this wave packet is that it represents a probability distribution. This means that where the amplitude of the wave packet is the greatest corresponds to the position in space with the highest probability of finding, or measuring, the particle. When the quantum wave packet is incident upon a barrier, it is partially reflected off the barrier and partially transmitted through the barrier. Since the packet transmitted through the barrier is a portion of the original probability distribution there is a small but finite probability of measuring the location of the quantum particle on the far side of the barrier. This phenomenon is known as tunneling and is well-known and accepted. However, a question arises as to the time required for the particle to achieve barrier penetration.

Several groups studying the phenomena of tunneling have shown that the tunneling velocities, or interaction times, for a variety of particles to pass through a barrier exceed  $c$ . For example, superluminal velocities have been measured for light pulses traveling through an absorbing material. Superluminal velocities have also been measured for the propagation for microwaves through a "forbidden zone" inside square metal waveguides. For a more detailed discussion of these experiments see, NEW SCIENTIST, vol. 146, pg. 27 (1995).

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More recently, a group at the University of California at Berkeley measured superluminal tunneling times for visible light tunneling through a dielectric mirror using a Hong-Ou-Mandel interferometer. Similar experiments by a group in the  
 5 University of Vienna in 1994 confirmed the Berkeley study and also showed that superluminal tunneling times could be obtained for increasingly large barrier thicknesses. For a more detailed discussion of these experiments see, NEW SCIENTIST, vol. 146, pg. 29 (1995).

Finally, in 1995, a group headed by Prof. Nimtz sent a microwave signal  
 10 broadcasting Mozart's 40<sup>th</sup> Symphony across 12 cm of space at 4.7 times the speed of light. For a more detailed discussion of this experiment see, NEW SCIENTIST, vol. 146, pg. 30 (1995).

In effect, these experiments show that tunnel times are independent of tunnel length, demonstrating the Hartman effect and tunneling. Under this regime the  
 15 tunneling time,  $\Delta t$ , is a saturated value and the Heisenberg uncertainty principle is written as follows:

$$\Delta \tau \Delta E = \hbar(1 + O) / 2 \quad (3)$$

20 where,  $\hbar$ , is the Heisenberg constant and,  $O$ , represent the higher order corrections to the tunneling time. This principle is referred to as the "energy borrowing" uncertainty principle, where the energy  $\Delta E$ , must be "paid back" in a time less than  $\Delta t$ , regardless of the energy flow speed or group velocity required to do so. A more detailed explanation of the physics of tunneling is provided in the following  
 25 references, each incorporated herein by reference: R.Y. Chiao, "Tunneling Times and Superluminality: a Tutorial", *quant-ph/9811019*, 7 Nov. 1998, at LANL; J. Jakiel et al., "On Superluminal Motions in Photon and Particle Tunnelings", *quant-ph/9810053*, 16 Oct. 1998, at LANL; A. Kempf, "A generalized Shannon Sampling Theorem, Fields at the Plank Scale as Bandlimited Signals", *hep-th/9905114*, 2  
 30 Mar. 2000, at LANL; P. Bamberg and S. Sternberg, "A course in Mathematics for Students of Physics 2", *Cambridge University Press 1990, Sect. 21.4*; J. Rembielinski, "Superluminal Phenomena and the Quantum Preferred Frame", *quant-ph/0010026*, 6 Oct. 2000, at LANL; J. Rancourt, "Optical Thin Films User Handbook", *SPIE Optical Engineering Press, 1996, Appendix C*; Hawking & Ellis,

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“The Large Scale Structure of Space-Time”, *Cambridge University Press, 1973, Sect. 4.3.*

While these experiments and texts clearly show the possibility of transmitting various forms of electromagnetic radiation faster than the speed of light, thus far no system has been developed to determine the one-way velocity vector of light utilizing these superluminal energy transmissions.

#### SUMMARY OF THE INVENTION

The present invention is directed to a device, system and method for measuring the one-way velocity of light using selective transmission technology to provide a superluminal energy flow.

This invention utilizes selective transmission technology to provide a superluminal energy flow. Selective transmission technology actively selects the high-energy components within a wavepacket for transmission through a Quantum barrier. The selective transmission technology or device accomplishes this by choosing a Quantum barrier, or air gap length, that selectively transmits only the high-energy portion of a wavepacket. These high-energy components are located near the wavepacket front or wavefront. The selective transmission device can then transmit these wavefront components more efficiently, giving these components a head start. These wavefront components contain Quantum information that is then used to completely reconstruct the wavepacket on the far side of the barrier before the energy of a free photon would have arrived on the far side. The transmitted wavefront information is used to completely reconstruct the wavepacket with energy borrowed from the vacuum on the far side of the barrier. The energy in the photon before the barrier must traverse the barrier at a speed that is FASTER than the vacuum speed of light. This is required to “pay-back” the energy borrowed from the vacuum on the far side of the barrier in a time that is equal to the time allowed by Quantum mechanics, or by the saturated Heisenberg energy borrowing uncertainty principle. This Quantum requirement along with the selective-transmission technology generates superluminal group velocities and superluminal energy flow. In summary, the chosen air gap length amplifies the front part, or high-energy components of the wavepacket using energy borrowed from the vacuum and Quantum information provided by selectively transmitted wavefront wavepacket components to completely reconstruct the wavepacket on the far side of the

Quantum barrier. This causes superluminal energy flow that is required to pay back the energy debt within the time required by Quantum mechanics. However, because of the time it takes to prepare the energy for superluminal transmission using selective-transmission technology the superluminal energy flow contains no superluminal classical-information.

If the single Preferred reference Frame (PF) has a small velocity,  $v(\text{PF})$ , relative to the selective transmission technology, then absolute causality requires that  $c(\text{AB}) - c \sim v(\text{PF})$ . The resulting vector tunnel time  $\Delta\tau(\text{AB})$  is proportional to  $1/(1 + \beta(\text{PF}))$ . It turns out that this preferred reference frame is also defined by the cosmic microwave background radiation Doppler shift, where at the Earth  $\beta(\text{PF}) = 0.001237 \pm 0.000002$ .

The Preferred reference Frame equivalence to zero Doppler shift in the cosmic microwave background is measurable. When the tunneling direction is in the direction of the red shift in the cosmic microwave background the tunneling time is shortest, and when the tunneling is in the blue shift direction the tunneling time is longest. In addition, the motion of the Earth around the sun adds and subtracts from the Earth's velocity relative to the cosmic microwave background preferred reference frame so that as the Earth rotates the tunneling direction, a vector tunneling time daily oscillation and yearly oscillation can be determined. Because the measured daily oscillation of the tunnel time is equivalent to the change in the vector vacuum velocity of light with tunneling direction and the tunneling direction is itself equivalent to the cosmic microwave background dipole direction created by the Doppler shift caused by the Earth's motion, the one-way light velocity can be measured.

In light of the above, in one embodiment, the invention is directed to a superluminal transmitter device comprising a transmission source, a receiver, and a selective-transmission device for receiving the transmission from the transmission source and selectively transmitting only the high-energy portion of the transmission wavefront through a barrier.

In a particular embodiment, the selective-transmission device comprises an air-gap barrier having proximal and distal ends formed from effective transmission barriers and an air gap disposed between the proximal and distal barriers such that a transmission from the transmission source enters the proximal end of the barrier tunnels across the air gap and exits the distal end of the barrier. In this

embodiment, the length of the air-gap is dependent on the wavelength of the wave-packet transmission such that the length of the air-gap is adjusted to efficiently transfer the higher energy part of the transmission wave-packet.

In another particular embodiment, the transmission source comprises a radio source in signal communication with a transmission antenna and the receiver comprises an amplifier in signal communication with a receiver antenna.

In yet another particular embodiment, the invention is directed to a speedometer and compass, comprising using the selective transmission device to measure the vector velocity of light relative to the absolute reference frame.

In still another particular embodiment, the invention is directed to a clock and calender, comprising using the selective transmission device to measure the vector velocity of light relative to the Earth's motion.

In still yet another particular embodiment, the invention is directed to a method for measuring the velocity vector of light. The method comprising measuring the oscillation in the tunneling time using a selective transmission system as described above.

In still yet another particular embodiment, the invention is directed to a method for calibrating temporal data. The method comprising measuring the oscillation in the tunneling time using a selective transmission system as described above.

## BRIEF DESCRIPTION OF THE DRAWINGS

These and other features and advantages of the present invention will be better understood by reference to the following detailed description when considered in conjunction with the accompanying drawings wherein:

FIG. 1 is a schematic view of an embodiment of the superluminal energy transmission device according to the invention.

FIG. 2 is a graphical representation of the of the superluminal transmission properties of the present invention.

FIG. 3 is a graphical representation of the superluminal transmission properties of the present invention.

FIG. 4 is a graphical representation of the superluminal transmission properties of the present invention.

FIG. 5 is a graphical representation of the superluminal transmission properties of the present invention.

FIG. 6 is a graphical representation of the superluminal transmission properties of the present invention.

FIG. 7 is a graphical representation of the superluminal transmission properties of the present invention.

FIG. 8 is a graphical representation of the superluminal transmission properties of the present invention.

## DETAILED DESCRIPTION OF THE INVENTION

The present invention is directed to a superluminal transmission device for measuring the vector velocity of light. In one embodiment, as shown in FIG. 1, the superluminal transmission device 10 comprises a transmission source 12, a selective-transmission device 14 adapted to receive a transmission from the transmission source 12, a receiver 16 in signal communication with the selective-transmission device 14 and a monitor 18 adapted to communicate the transmission to a user.

A transmission wavepacket 20 having a wavefront component and an energy component is introduced into the selective-transmission device 14 from the transmission source 12 such that the transmission wavepacket 20 is conducted through the space between the transmission source 12 and the receiver 16 to the monitor 18 at velocities faster than the speed of light. The selective-transmission device 14 is placed in proximate relation to the transmission source 12 such that the transmission wavepacket 20 passes through the selective-transmission device 14 and the wavefront component of the transmission wavepacket 20 is transmitted into the receiver 16 creating a signal. A receiver or series of receivers 16, are adapted to receive the signal and transmit the signal to a monitor 18 in signal communication therewith. Any device having the ability to detect changes in amplitude, frequency, phase or wavelength of the transmission 20 can be used as a receiver 16 and monitor 18, such as, for example, a radio amplifier in signal communication with an oscilloscope or a Time to Digital Converter (TDC). Additionally, any suitable transmission source 12 may be used in the subject invention, such as, for example, a optical laser, a microwave generator or a radio

transmitter so long as detectable levels of electromagnetic radiation are transmitted to the receiver 16 in the form of a transmission wavepacket 20.

In general terms, the selective-transmission device 14 comprises a quantum air-gap barrier 22, which is in signal communication with the transmission source 12. The quantum air-gap barrier 22 comprises a proximal 24 and distal 26 barrier wall and an air-gap 28 having a tunneling, or air-gap, length 30 disposed therebetween. The proximal barrier wall 24 is in signal communication with the transmission source 12 and the distal barrier wall 26 of the air-gap barrier 22 is in signal communication with the receiver 16. The transmission 20 from the transmission source 12 interacts with the air-gap barrier 22 which selectively transmits the high-energy component of the transmission wavepacket 20 across the air-gap 28 to the receiver 16 at superluminal velocities. The air-gap barrier 22 generates superluminal transmission velocities in the high-energy component of the transmission 20 by selecting the high energy component of the transmission wavepacket 20 and more efficiently transmitting that high-energy component across the air-gap 28. The high-energy component of the transmission wavepacket 20, is selected by arranging the proximal 24 and distal 26 barrier walls such that the air-gap length 30 therebetween corresponds to quarter wavelength or multiples thereof of the higher energy component of the transmission wavepacket 20. By selecting the air-gap length 30 to correspond to the wavelength of the high energy component of the total transmission wavepacket 20, the air-gap barrier 22 provides the high-energy component of the transmission wavepacket 20 a head start, in effect causing tunneling of the high-energy component, or tunneling transmission across the air-gap 28 in a tunneling time that is independent of the tunnel distance, or air-gap length, 30 according to the Hartman effect, thus causing the tunneling transmission to cross the air-gap 28 at a superluminal velocity. Any air-gap barrier 22 construct suitable for selecting the high energy component of the transmission wavepacket 20 from a transmission source 12 and transmitting the high-energy component of the wavepacket 20 across an air-gap 28 at superluminal velocities may be used such as, for example, optical mirrors for visible light transmissions, square metal waveguides for microwave transmissions or tanks of having a high index of refraction substance such as water for radio transmissions.

In one preferred embodiment, a radio transmission source 12, a radio receiver 16 and an air-gap barrier 22 comprising a proximal tank 24 and a distal



1 tank **26** aligned parallel to each other across an air-gap **28** are utilized to generate  
the superluminal transmissions. The proximal tank **24** is placed in signal  
communication with the transmission source **12** and the distal tank **26** is placed in  
5 signal communication with the receiver **16**. The tanks **24** and **26** are arranged such  
that an air-gap **28** is created between having an air-gap length **30** corresponding to  
a quarter wavelength of the high-energy component of the transmission  
wavepacket **20**. In this embodiment, the tanks **24** and **26** may have any index of  
refraction suitable to act as a quantum barrier such as, for example, a plexiglass  
10 tank filled with water.

To transmit the transmission wavepacket **20** to and from the selective-  
transmission device **14**, the transmission source **12** and receiver **16** must be  
positioned relative to selective-transmission device **14** such that the transmission  
wavepacket **20** passes through the selective-transmission device **14**. In the  
15 embodiment shown in the attached figures, a radio transmission source **12** and a  
radio receiver **16** utilize antennas **32** directed at the selective-transmission  
device **14**. However, any suitable design can be used such that the transmission **20**  
from the transmission source **12** passes through the selective-transmission  
device **14** and enters the receiver **16**. For example, focusing optics can be used to  
20 direct a visible light transmission **20** from a laser transmission source **12** through  
a selective-transmission device **14** to a photomultiplier receiver **16**.

A prototype of the superluminal transmission device **10** described above was  
constructed. A NIM-logic pulser **34** (Phillips Scientific model 417 Nuclear  
Instrumentation Standard Pocket Pulser) in signal communication with an  
25 amplifier **36** (RadioShack catalog# 15-1113C) is used as the transmission source **12**  
and is placed in signal communication with a five-element folded-dipole Yagi  
antenna **32a** designed for two-meter wavelength radio waves. A second amplifier **38**  
(RadioShack catalog# 15-1170) in signal communication with a second five-element  
folded-dipole Yagi antenna **32b** is used as the receiver **16**. Both antennas **32a**  
30 and **32b** comprise ¼ inch aluminum ground wire reflector and deflectors, and a #10  
copper wire folded dipole. 75 ohm to 300 ohm transformers, (RadioShack catalog  
# 15-1140), are connected to 75 ohm cables at the antennas **32a** and **32b**. Each  
antenna **32a** and **32b** is also surrounded by an aluminum screen (not shown), with  
a 114 cm wide opening along the folded-dipole direction to selectively transmit and  
35 receive a signal wavelength at  $\leq 228$  cm. The signal from the receiver amplifier **16**

is fed into an oscilloscope monitor **18** (Tektronix TDS220). Alternatively a TDC could be utilized as a monitor **18**, such as, for example, an ORTEC 9308 Picosecond Time Analyzer preceded by a 9307 pico-Timing Discriminator. The transmission source **12** signal is also monitored by the oscilloscope monitor **18** via a signal splitter **40** which is placed in signal communication with the radio-wave pulser **34**. The cables leading from the transmission source **12** and the receiver **16** to the oscilloscope monitor **18** are terminated into 75 ohms.

The selective-transmission device **14** comprises an air-gap barrier **22** having proximal **24** and distal **26** barrier walls arranged such that an air-gap **28** lies therebetween. The proximal **24** and distal **26** barrier walls consist of two 4 ft wide and 2 ft high distilled water tanks. The distilled water layer thickness in each tank is 12.7 mm or  $\frac{1}{2}$  inch and the index of refraction is  $n = 9$  and  $k = 0.002$ . The water tanks are constructed with  $\frac{1}{4}$  inch thick Plexiglass having an index of refraction of  $n = 1.6$  and  $k = 0.0$ . The proximal **24** and distal **26** barrier walls can be adjusted such that the air-gap length **30** between them extends up to 270 cm.

FIGs. 2 to 7 show the results of a typical superluminal transmission absent a signal pulse for the superluminal transmission device prototype **10** shown in FIG. 1. During a transmission, the source amplifier **36** gain is set at the minimum level and the FM trap is turned off. The cable lengths are adjusted such that the pulser **34** trigger pulse arrives at the oscilloscope monitor **18** just prior to the transmission wavepacket wavefronts **20**. Each transmission measurement contains 128 samples, averaged by the oscilloscope monitor **18**. The source data, or standard is taken with only the proximal barrier wall **24** in place. All error bars are the standard deviation of five data set measurements. FIG. 2 shows data from a source wavepacket measurement. The measured peak to peak time,  $\tau_m$ , for the source wavepacket is  $7.6 \pm 0.1$  ns, giving a photon wavelength of 228 cm. The large pulse shown below 0 peaking time,  $\tau_p$ , is the pulser trigger which is the rising edge set at -0.4 volts. The peaking time of the source wavepacket,  $\tau_p$ , relative to the pulser trigger is  $39.0 \pm 0.3$  ns. Table 1, below, shows the peak to peak separation times of the various components of the transmission wavepacket **20**. As shown for peak numbers 1 to 3, the higher energy, or lower wavelength, components are in the front part, or nearer the wavefront, of the transmission wavepacket **20**. It is these high energy components of the transmission wavepacket **20** that are selectively transmitted by the superluminal transmission device **10** described above.

Table 1: Peak to Peak Separation Times		
Peak Numbers	Peak to Peak Time (ns)	Wavelength (cm)
1 to 2	6.8	204
2 to 3	6.8	204
3 to 4	7.2	216
4 to 5	$7.6 \pm 0.1$	228
5 to 6	7.6	228

In FIG. 3, the transmitted energy,  $E_T$  ( $\text{mV}^2$ ), averaged over time,  $\langle E_T \rangle$ , from 0 to 80 ns, is shown versus the air-gap length,  $L$  (cm). The maximum transmitted energy,  $E_T$ , over 0 to 80 ns occurs at an air-gap length of 57 cm. This energy peak is identified as a 228 cm photon quarter wavelength, indicating that the shorter air-gap lengths transmit the quarter wavelength or higher energy components of the transmission wavepacket **20** more efficiently. This is the mechanism that preferentially selects the front part, or high energy components, of the wavepacket **20** for transmission and generates the superluminal group velocities shown for a 30 cm air-gap length **30**. The tunneling time,  $\Delta\tau$  (ns), is also shown verse the air-gap length,  $L$  (cm), in FIG. 3. The flat tops of the shaded boxes identify the regions where the tunneling time is independent of the tunnel, or air-gap, length **30**. This is a demonstration of the Hartman effect discussed above. As previously described, the tunneling time is defined by the “energy borrowing” Heisenberg uncertainty principle of Equation (3), where the energy must be “paid back” in a time less than the tunneling time regardless of the energy flow speed or group velocity required. The wavepacket **20** that tunnels through the air-gap **28**, peaks prior to the source, or non-tunneling, wavepacket. The measured peaking time difference is defined by the equation:

$$\tau_g = \tau_p - \tau_{psource} \quad (4)$$

where,  $\tau_g$ , is the measured group delay time. The tunnel time is then defined by the equation:

$$\Delta \tau = (L/c) + \tau_g \quad (5)$$

where,  $(L/c)$ , is the source time. Peaking times, group delay times and tunnel times as measured during a transmission measurement are listed in Table 2, below.

Table 2: Wavepacket Peak Times, Group Delay Times and Tunnel Times			
Air-Gap Length (cm)	$\tau_p$ (ns)	$\tau_g$ (ns)	$\tau$ (ns)
source	$38.96 \pm 0.33$	—	—
200	$38.05 \pm 0.26$	$-0.91 \pm 0.42$	$5.75 \pm 0.42$
210	$37.69 \pm 0.33$	$-1.27 \pm 0.47$	$5.73 \pm 0.47$
220	$37.35 \pm 0.26$	$-1.60 \pm 0.52$	$5.73 \pm 0.52$
230	$37.07 \pm 0.48$	$-1.89 \pm 0.58$	$5.77 \pm 0.58$
240	$36.77 \pm 0.60$	$-2.19 \pm 0.68$	$5.80 \pm 0.68$

As shown, the tunnel time is independent of the length of the air-gap **28**, indicating an increase in the negative group-delay time for the tunneling portion of the transmission wavepacket **20** as the air-gap length **30** is increased. As shown by FIG. 3, for air-gap lengths between 200 and 240 cm the tunnel time is less than the source time or vacuum speed of light,  $c$ . The increase in the tunnel time standard deviations with increasing tunnel length measure a lower bound in the tunnel time distributions that are proportional to the tunnel length.

FIG. 4, shows a graph plotting the transmission fraction and tunneling time for a 204 cm wavelength photon verses the air-gap length. The boxes in the figure again show the regions where the tunnel time is independent of the tunnel length. This plot also shows that those regions where the tunnel time is independent of the tunnel length coincide with the selective transmission of the wavepacket front, the 204 cm, or high-energy, wavepacket components. Thus, selective transmission of the high-energy component of the transmission wavepacket **20** causes superluminal group velocities in those transmitted components.

FIG. 5 graphically shows the effect of this selective transmission. Data for this graph was taken from one 128 sample source data set for an air-gap length of 220 cm. In this measurement a 204 cm wavelength wavepacket **20** with a transmission fraction of 0.997 is utilized and the wavepacket front is being actively

selected. The straight bold line shows the vacuum speed of light, while the bent bold line depicts the speed of the selective transmission of the high-energy component of the wavepacket **20**. As shown, the selective transmission creates velocities in the high-energy components that are much less than  $c$  but not enough less to hold the energy luminal. In mathematical terms, the stationary phase tunnel time,  $\tau_s$ , given by:

$$\tau_s = \partial\phi / \partial\omega \quad (6)$$

added to the head-start time,  $\tau_h$ , caused by selecting the high-energy wavepacket components equals the tunnel time,  $\Delta\tau$ . The stationary phase tunnel time, give by Equation 8, is  $\tau_s = 46.18$  ns, its peak value. The head-start time is defined by the equation:

$$\tau_h = \Delta\tau - \tau_s \quad (7)$$

where,  $\tau = 5.73$  ns, and  $\tau_h = -40.45$  ns. Under these conditions the head-start time is also defined by the equation:

$$\tau_h = \tau_g - \tau_{psource} \quad (8)$$

where,  $\tau_h = -40.56$  ns. Indicating that by selecting these high energy components **42** a 40.56 ns head-start can be generated in the transmission, advancing the energy 1.6 ns for superluminal energy flow.

FIGs. 6 and 7 show tunnel and source data for single data sets, each containing 128 samples averaged by the scope. The tunnel data for FIG. 6 is from a measurement taken utilizing an air-gap length of 220 cm. The negative group delay time of -2.1 ns is for this data set. The group delay time average over all five data seta taken at 220 cm is -1.6 ns. In this plot, the tunneled wavepacket contains less energy than the source wavepacket. However, from the transmission fraction analysis it is shown that the front part of the wavepacket is amplified more than the tail part, producing negative group delay times. For the 30 cm air-gap length, shown in FIG. 7, the tunneled energy is greater than the source energy as shown in FIG. 3 and the tunneled wavepacket hugs the trailing edge of the source

1 wavepacket. This effect identifies a causality restriction. The tunneled and source  
 wavefronts must be simultaneous as required by the special theory of relativity.  
 Causality requires the computed tunnel time to be  $\tau = [(2L/c) - \tau_s]$  for air-gap lengths  
 5 of 30 through 70 cm.

FIG. 8 shows a measurement of the daily oscillation of the tunneltime, which  
 is equivalent to the change in the vector vacuum velocity of light with tunneling  
 direction. This tunneling direction is in turn equivalent to the cosmic background  
 dipole direction created by the Doppler shift caused by the Earth's motion.  
 10 Accordingly, the one way light velocity can be measured.

As a short explanation, the single Preferred reference Frame (PF) has a small  
 velocity,  $v(\text{PF})$ , relative to the selective transmission technology, then absolute  
 causality requires that  $c(\text{AB}) - c \sim v(\text{PF})$ . The resulting vector tunnel time  $\Delta\tau(\text{AB})$   
 is proportional to  $1/(1 + \beta(\text{PF}))$ . It turns out that this preferred reference frame is  
 15 also defined by the cosmic microwave background radiation Doppler shift, where at  
 the Earth  $\beta(\text{PF}) = 0.001237 \pm 0.000002$ . The vector  $\beta(\text{PF})$  points in the direction of  
 declination  $7.22^\circ$  closest to the sun on March 7<sup>th</sup> with a right ascension of 23.20 h  
 and is in the opposite direction to the Earth's motion that causes the cosmic  
 microwave background Doppler shift.

20 The Preferred reference Frame equivalence to zero Doppler shift in the  
 cosmic microwave background is measurable. The theoretical vector tunneling time  
 utilizing this selective transmission technology is  $\Delta\tau(\text{AB}) = \Delta\tau \pm 7.08 \pm 0.57$  ps  
 where the Doppler shift defines the  $\pm 7.08$  ps of daily oscillation and the  $\pm 0.57$  ps  
 is yearly oscillation and is a measure of the Earth's motion around the sun.

25 The measured daily oscillation of the tunnel time is equivalent to the change  
 in the vector vacuum velocity of light with tunneling direction. When the tunneling  
 direction is in the direction of the red shift in the cosmic microwave background the  
 tunneling time is shortest or plus 7.08 ps. When tunneling is in the blue shift  
 direction the tunneling time is longest or plus 7.08 ps.

30 The measured daily oscillation of the tunnel time is due to a change in the  
 vector vacuum velocity of light  $c(\text{AB})$ , as a function of tunneling direction (AB), and  
 the Reichenbach clock coefficients, as described in Equations (1) and (2). Utilizing  
 the prototype system, the one-way light velocity can be measured and compared with  
 these theoretical values.

It is found that measurements taken with the prototype track these theoretical values. The vector vacuum velocity of light was measured at an air-gap length of 220 cm. FIG. 8 shows a histogram mean value data of the tunneling time over a twenty-four hour period of measurement. The 9308 has a histogramming bin width of 1.22 ps over the 80 ns window. At  $L = 220$  cm, the standard deviation lower bound,  $\Delta\tau(\min)1 = \Delta X/2c = 507\text{ps}$ , requiring millions of pulser pulses to decrease the error in the tunneling time histogram mean value below a picosecond. In order to maximize the group delay time the 9307-discriminator level was set as high as possible without effecting the count rate. The tunneling direction was parallel to the Earth's surface at  $108^\circ$ , fixing the tunneling direction declination at  $-12^\circ$ . A typical data set showing peak time statistic in ns for ten spectrum centroids is summarized in Table 3, below:

Table 3: Tunneling Time Oscillation Data	
Mean	47.010
Median	47.011
RMS	47.010
Standard Deviation	0.0029258
Variance	$8.5604 \text{ e}^{-6}$
Standard Error	0.00092523
Skewness	1.0718
Kurtosis	0.97325

As described above, the cosmic Doppler shift defines the theory line shown in FIG.8. As the measured daily oscillation of the tunnel time is equivalent to the change in the vector vacuum velocity of light with tunneling direction, and as the tunneling direction is equivalent to the cosmic microwave background dipole direction created by the Doppler shift caused by the Earth's motion, the one way light velocity is measured, and the Reichenbach clock synchronization coefficient determined as:  $\beta(\text{QH}) = 0.001237 \pm 0.000002$  in the 23.20h right ascension and  $7.22^\circ$  declination direction in the Doppler red-shift direction.

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Although the above embodiment was only utilized to measure the one way light velocity, it will be obvious to one of skill in the art that other uses for the oscillating tunneling time measurements could be made. For example, using the tunneling time oscillation to calibrate data in the time domain. It is clear that by knowing that a predictable and computable oscillation exists in temporal data, a calibration factor can be determined and utilized to calibrate that temporal data utilizing the invention. In addition, the vector velocity of light can be used as a speedometer and compass relative to the absolute reference frame or as a clock and calender by also knowing the Earth's motion.

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In addition, the measured preferred reference frame equivalence to the cosmic microwave background preferred reference frame is very aesthetic. This measured equivalence begins the new science of "Radio Cosmology". It will be understood that as the special theory of relativity becomes more complicated to accommodate absolute causality, Minkowski space-time gives way to, the "more natural" Lorentz frame bundle with its base space of preferred frame velocities. This "more natural" frame bundle language should be explored using more accurate "Radio Cosmology" measurements. For example, more accurate measurements are required to confirm Lorentz over Galilean symmetry using only tunneling time data

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Although specific embodiments are disclosed herein, it is expected that persons skilled in the art can and will design alternative light velocity vector measurement systems that are within the scope of the following claims either literally or under the Doctrine of Equivalents.

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